

Lecture 3

String interactions

closed, oriented

$$A(1, \dots, M) = \begin{array}{c} \text{xxx} \\ \bigcirc \end{array} + \begin{array}{c} \text{xxx} \\ \bigcirc \\ \text{a} \end{array} + \begin{array}{c} \text{a} \quad \text{a} \\ \text{a} \quad \text{a} \end{array} + \dots$$

$$= \sum_{g=0}^{\infty} \underbrace{\int \mathcal{D}X \mathcal{D}h e^{-S_P[X, h]} V_1 \dots V_M}_{\text{fixed topology}}$$

$M=0$

$$Z = Z_g = \int \mathcal{D}h \mathcal{D}X e^{-S_P[X, h]}$$

$$= \bigcirc + \bigcirc + \begin{array}{c} \text{a} \quad \text{a} \\ \text{a} \quad \text{a} \end{array} + \dots$$

fixed topology

no external states

closed, oriented strip/surfaces

gauge fixing: $h = e^{2\phi} \overset{\uparrow}{h}$
 locally $\overset{\uparrow}{h} = \begin{pmatrix} \eta & \text{Lor} \\ S & \text{Em} \end{pmatrix}$

Consider PE for fixed metric

$$\int DX e^{-S_P[X, h]}$$

Gaussian integrals

$$\int_{-\infty}^{\infty} dx e^{-ax^2 + bx + c} = \sqrt{\frac{\pi}{a}} e^{\frac{b^2}{4a} + c}$$

$$Re(a) > 0$$

$$A \in GL(n, \mathbb{R}) \quad A = A^T$$

$$\int d^n x e^{-\frac{1}{2} x^T A x + b x} = \sqrt{\frac{(2\pi)^n}{\det A}} e^{\frac{1}{2} b^T A^{-1} b}$$

$$\begin{aligned} S[\phi] &= \frac{1}{2} \int d^n x \sqrt{h} (h^{mn} \partial_m \phi \partial_n \phi + m^2 \phi^2) \\ &= \frac{1}{2} \int d^n x \sqrt{h} \phi (\Delta_h + m^2) \phi \geq 0 \end{aligned}$$

$$\Delta_h = \Theta \frac{1}{\sqrt{h}} \partial_m \sqrt{h} \cdot h^{mn} \partial_n$$

$$\int \mathcal{D}\phi e^{-S[\phi]} = \det^{-1/2} (\Delta_{hT} m^2)$$

$m > 0 \rightarrow$ fermion = two modes (units)

ζ -det's

$$\zeta(s) = \sum_{n=1}^{\infty} n^{-s} \quad \text{Re } s > 1$$

extends merom to $\mathbb{C}$, simple pole at $s=1$

otherwise holomorphic

$$\zeta(-1) = -\frac{1}{12}, \quad \zeta(0) = -\frac{1}{2}, \quad \zeta'(0) = -\frac{1}{2} \ln(2\pi)$$

applies

$$\left(\sum_{n=1}^{\infty} n \right)_{\zeta} = \zeta(-1) = -\frac{1}{12}$$

$$\left(\sum_{n=1}^{\infty} 1 \right)_{\zeta} = \zeta(0) = -\frac{1}{2}$$

$$\left(\prod_{n=1}^{\infty} a \right)_r = a \left(\sum_{n=1}^{\infty} 1 \right)_r = a^{-\frac{1}{2}}$$

$$A = A^+ \quad / \quad \lambda_1 \dots \lambda_N > 0$$

$$\zeta_A(s) = \sum_{n=1}^N \lambda_n^{-s} \quad s \in \mathbb{C}$$

converges for large enough $\operatorname{Re} s$

$$e^{-\zeta'_A(0)} = \det A$$

$A =$ selfadj non-neg diff operator
w/ discrete spectrum $\{\lambda_n\}$ on M ,
compact

$$\zeta_A(s) = \sum_n \lambda_n^{-s} \quad s \in \mathbb{C}$$

omit zero modes

conv. for $\operatorname{Re} s > \frac{\dim M}{\deg A}$
regular at $s=0$

$$\det' A = e^{-\zeta'_A(0)}$$

compute

$$Z = N \int \mathcal{D}X e^{-S_P[X; h]}$$

$$S_P[X; h] = \frac{1}{2} \int_{\Sigma} d^2z \sqrt{h} h^{\alpha\beta} \partial_{\alpha} X^{\mu} \partial_{\beta} X_{\mu}$$

$$= \frac{1}{2} \int_{\Sigma} d^2z \sqrt{h} X^{\mu} \Delta_h X_{\mu}$$

$$= \frac{1}{2} (X, \Delta_h X)$$

$$\Delta_h = - \frac{1}{\sqrt{h}} \partial_{\alpha} \sqrt{h} h^{\alpha\beta} \partial_{\beta}$$

$$\Delta_h \psi_m = \lambda_m \psi_m \quad (\langle \psi_m | \psi_n \rangle = \delta_{m,n})$$

$$X^{\mu} = \sum_{n=0}^{\infty} a_n^{\mu} \psi_n = X_0^{\mu} + X^{\mu}$$

$$a_n^{\mu} \in \mathbb{R}$$

$$1 = (\psi_0, \psi) = \int_{\Sigma} d^2z \sqrt{h} \psi_0^2$$

$$\Rightarrow \psi_0 = \left(\int_{\Sigma} d^2z \sqrt{h} \right)^{-1/2} = \text{vol}_h(\Sigma)^{-1/2}$$

$$X_0^{\mu} = a_0^{\mu} \psi_0$$

$$Z = N \int DX e^{-\frac{1}{2} (X, \Delta_n X)}$$

$$= N \int \prod_{m=1}^M da_m^M e^{-\frac{1}{2} \sum_{n=1}^N \lambda_n (a_n^M)^2}$$

$$= N \left(\int \prod_{\mu} da_{\mu}^0 \right) \left(\int \prod_{n=1}^N \prod_{m=1}^M da_m^n e^{-\frac{1}{2} \sum_{n=1}^N \lambda_n (a_n^M)^2} \right)$$

$$\int \prod_{\mu} da_{\mu}^0 = \text{vol}(M)$$

$$X_0^M = a_0^M \gamma_0$$

$$\int \prod_{\mu} da_{\mu}^0 = \left(\int \prod_{\mu} dX_0^{\mu} \right) \gamma_0^{-D}$$

$$= \text{vol}(M) / \text{vol}_n(\mathbb{R})^{D/2}$$

$$\int \prod_{\mu} \prod_{n=1}^N da_m^n e^{-\frac{1}{2} \sum_{n=1}^N \lambda_n (a_n^M)^2}$$

$$= \prod_{\mu} \prod_{n=1}^N \int \lambda_n^{-\frac{1}{2}} = \prod_{\mu} \lambda_n^{-D/2}$$

$$= (\det_g \Delta_n)^{-D/2}$$

$$Z = N \det M \left(\frac{\det g \Delta q}{\det_n \bar{\epsilon}} \right)^{-\frac{D}{2}}$$

Fermionic Gaussian integrals

$$\theta_i \theta^2 = 0$$

$$\int d\theta \theta = 1 \quad \int d\theta = 0$$

2n variables $\theta_i, \bar{\theta}_i \quad i=1, \dots, n$

$$\int d\theta_i \theta_j = \delta_{ij} = \int d\bar{\theta}_i \bar{\theta}_j$$

$$\int d\theta_i = 0 = \int d\bar{\theta}_i$$

Int. over all variables, sign convention

$$\int \prod d\theta d\bar{\theta} \prod_{i=1}^n (\bar{\theta}_i \theta_i) = 1$$

+

$$Z = \int D\theta D\bar{\theta} e^{-\bar{\theta}^T A \theta} = \det A$$

$$\text{writing } \det A = \sum_{\sigma \in S_n} \sum_{i=1}^n A_{i\sigma(i)}$$

~~Dirac gamma, matrices, flat Lorentz metric~~

$$\int D\psi D\bar{\psi} e^{-\int d^4x \bar{\psi} i \partial^\mu \gamma_\mu \psi}$$

$$= \det(i \partial^\mu \gamma_\mu) = \det(-\square)$$

$$\begin{aligned} i \partial^\mu \gamma_\mu i \partial^\mu \gamma_\mu &= -\{\gamma^\mu, \gamma^\mu\} \partial_\mu \partial_\mu \\ &= -\eta^{\mu\mu} \partial_\mu \partial_\mu = -\square \end{aligned}$$

Massless Dirac field, Euclidean $\mathbb{R}^m$

$$Z = \int \bar{\psi} \psi \, d\psi \, e^{-\int d^m x \, \bar{\psi} i \not{\partial} \psi}$$

$$= \det(i \not{\partial}) = (\det \Delta_S)^{\frac{1}{2}}$$

$$i \not{\partial} i \not{\partial} = - \gamma^\mu \partial_\mu \gamma^\nu \partial_\nu$$

$$= - \frac{1}{2} \{ \gamma^\mu \gamma^\nu \} \partial_\mu \partial_\nu$$

$$= - \delta^{\mu\nu} \partial_\mu \partial_\nu \equiv \Delta_S$$

Integration over moduli

$$Z = N \int D h D X e^{-S_P I(h)}$$

Sym: $P_{\alpha\beta} \propto W_{\alpha\beta}$

gauge fix $h_{\alpha\beta} = e^{2\phi} \tilde{h}_{\alpha\beta}$

locally $\tilde{h}_{\alpha\beta} = (g_{\alpha\beta}/i) \delta_{\alpha\beta}$

Further integration over moduli

$$\delta h_{\alpha\beta} = - \underbrace{(\nabla_\alpha v_\beta + \nabla_\beta v_\alpha)}_{\text{Per}} + \underbrace{2\Lambda h_{\alpha\beta}}_{\text{Weyl}}$$

$$= - \underbrace{(P_v)_{\alpha\beta}}_{\text{trunk}} + \underbrace{2\tilde{\Lambda} h_{\alpha\beta}}_{\text{trace}}$$

$$P: v_\alpha \rightarrow (P_v)_\alpha = \nabla_\alpha v_\beta + \nabla_\beta v_\alpha - \nabla_\gamma v^\gamma h_{\alpha\beta}$$

$$P^\dagger: t_{\alpha\beta} \rightarrow (P^\dagger t)_\alpha = -2\nabla^\beta t_{\alpha\beta}$$

$$t_{\alpha\beta} = t_{\beta\alpha}$$

$$h^{\alpha\beta} t_{\alpha\beta} \geq 0$$

$$(v, P^+ t) = (P_v, t)$$

$$D_h = D(P_v) D \tilde{\Lambda}$$

$$= D_v D \Lambda \underbrace{|dA P|}_{= (dA P P^+)^{1/2}}$$

assume: measure is Weyl invariant
in zero modes of P, P^+

vector fts $\xleftrightarrow{r=1}$ STT,
rt det of $\dagger$ vectors
rest of measure

$$Z = \underbrace{N(\int Dv) (\int D\Lambda)}_{=1} \int DX |dA P| e^{-S[X, e^{2h} \vec{b}]}$$

simultaneous matrices

$$\det M^T = \det M$$

$$\Rightarrow \det M^T = \det M^* = (\det M)^*$$

$$\begin{aligned} \Rightarrow \det M M^T &= \det M \det M^T \\ &= \det M (\det M)^* \\ &= |\det M|^2 \end{aligned}$$

$$\Rightarrow (\det M M^T)^{1/2} = |\det M|$$

Express FP determinant as
Gaussian integral over
continuous tensor field

$$|dA P| = \int D b D c e^{-\frac{1}{2\pi} \int d^2 x \left(\partial_\alpha c \partial_\alpha c + \sqrt{h} h^{\alpha\beta} b_{\alpha\beta} \nabla_\alpha c \right)}$$

$$= \int D b D c e^{-S_{\text{dual}}[b, c, h]}$$

$b_{\alpha\beta} = b_{\beta\alpha}$ $h^{\alpha\beta} b_{\alpha\beta} = 0$ *antisym*
 c^α *ghost* $\Leftrightarrow$ *xy* $\downarrow$
 $\uparrow$ *number*
 diff of
 metric

$$Z = \int D b D c D X e^{-S_P[X, h]} \rightarrow Z_h[b, c, h]$$

$$\boxed{\begin{matrix} X\text{-CFT} \\ c \end{matrix}} + \boxed{\begin{matrix} b, c\text{-CFT} \\ c_{\text{ghost}} \end{matrix}}$$

$\uparrow$
 needed for quantum gauge
 invariance

Meaning i.g. not Weyl invariant

$$Z = \int D\phi D\bar{\phi} D\psi D\chi e^{-S_\phi - S_{gh} - (C + g_h)\Sigma}$$

$$\Sigma = \Sigma[\phi, \bar{\psi}] \quad \text{Liouville action}$$

ϕ = Liouville field $\hat{=}$ trace part
of metric / w/ factor

• critical string

$$\left. \begin{array}{l} C + C_{gh} = 0 \\ C_{gh} = -26 \end{array} \right\} \Rightarrow D = L = 26$$

w/ anomaly cancels

• non crit. Str.

WS CFT acquires add dof ϕ

$D \neq 26$, complicated dynamics

zero modes

$$(P \psi)_\alpha = 0 \Rightarrow \psi_\alpha \text{ CKV}$$

$\hat{=}$ system which do not change
moduli

$\rightarrow$ overcounting in PI

$\rightarrow$ add gauge fixing

$\rightarrow$ separate vol (GC)

$$(P^\dagger h)_{\alpha\beta} = 0 \Rightarrow h_{\alpha\beta} \text{ moduli diff}$$

transverse moduli deform which
cannot be gen by vector fields

"undercounting", no global

gauge fixing

$$h_{\alpha\beta} > e^{2\phi} \hat{h}_{\alpha\beta}$$

On Σ_g :

$$h_{df} \approx \sum_{conf} e^{i \int h_{df}(m_i)}$$

matrix of conf interaction

$$\begin{aligned} M_g &= \{ \text{conf conf matrix on } \Sigma_g \} \\ &= \{ \text{conf starts \& on } \Sigma_g \} \\ &= \{ \text{cplx starts on } \Sigma_g \} \end{aligned}$$

$$\text{dim } M_g = \begin{pmatrix} 0 & g=0 \\ 1 & g=1 \\ 3g-3 & g \geq 1 \end{pmatrix}$$

- To count all moduli must
- + include integration over M_g

As

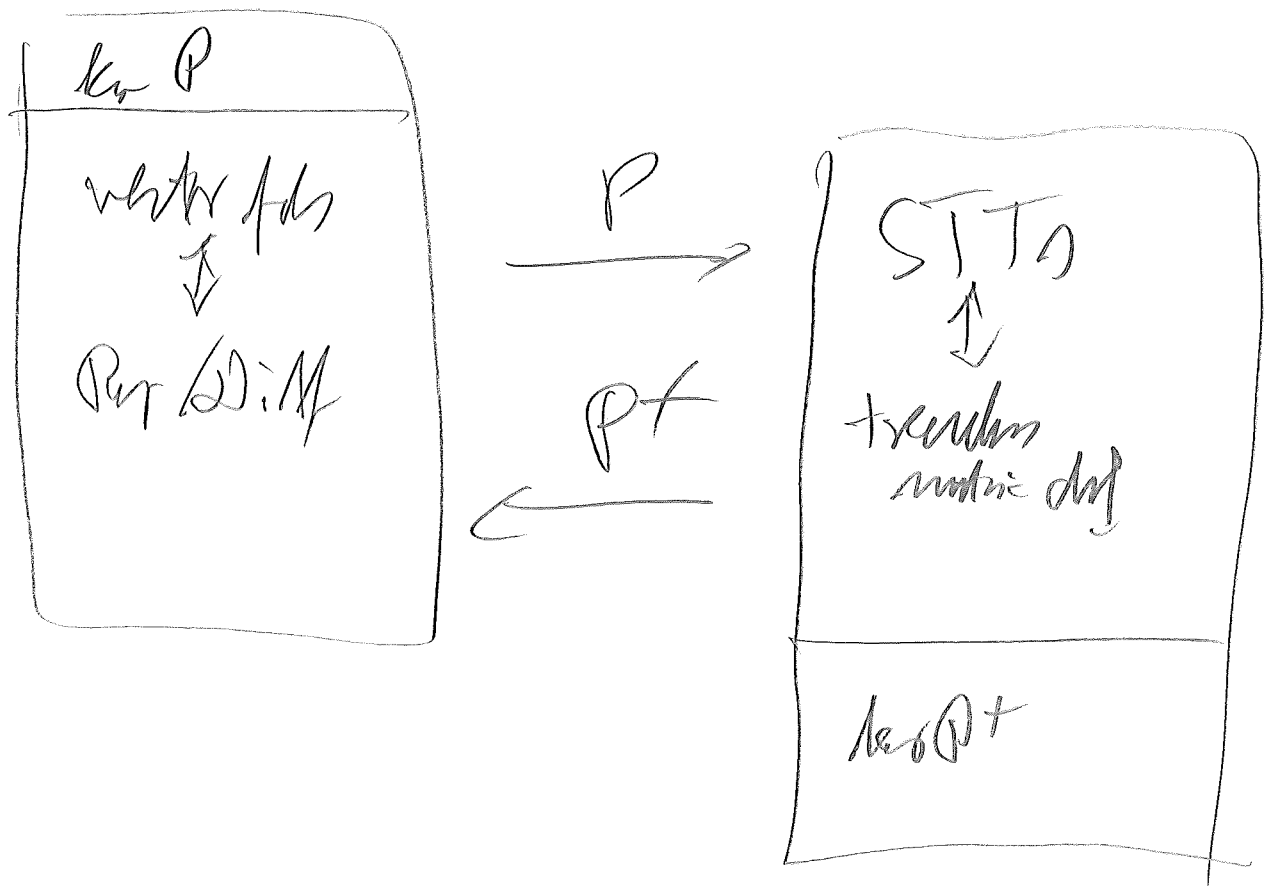
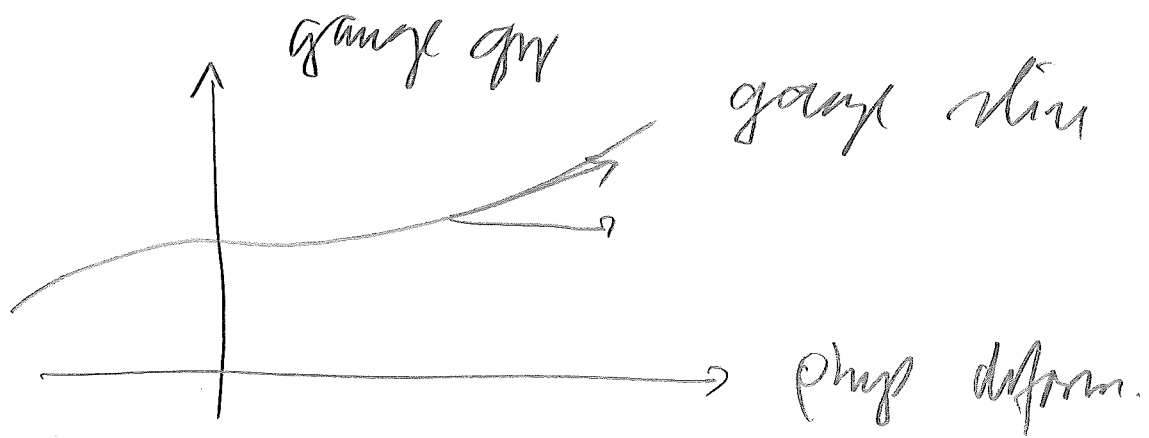
$$Z = N \text{Vol}(\mathcal{BC}) \int \prod_b D\phi D\psi D\chi \int_{M_g} d\mu(m_i) \\ e^{-\left(\int \bar{\psi} X_1 \psi(m_i) \right) - \int_{gh} \bar{\psi} X_1 \psi(m_i))}$$

CFT

$X, b, c,$

$on (Z_g, \mathbb{H}(m_i))$

$$c_{total} = c_x + c_{gh} = 26 - 26 = 0$$



$$v_l(\text{CKV}) (d\alpha'_g P^+)^{\frac{1}{2}} d_\mu(m_1 \dots)$$

A